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Flow-aligned jets in the magnetospheric cusp: Results from the Geospace Environment Modeling Pilot program

L. A. Weiss,^{1,2} P. H. Reiff,¹ E. J. Weber,³ H. C. Carlson,³ M. Lockwood,⁴
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Abstract. The extended flight of the Airborne Ionospheric Observatory during the Geospace Environment Modeling (GEM) Pilot program on January 16, 1990, allowed continuous all-sky monitoring of the two-dimensional ionospheric footprint of the northward interplanetary magnetic field (IMF) cusp in several wavelengths. Especially important in determining the locus of magnetosheath electron precipitation was the 630.0-nm red line emission. The most striking morphological change in the images was the transient appearance of zonally elongated regions of enhanced 630.0-nm emission which resembled "rays" emanating from the centroid of the precipitation. The appearance of these rays was strongly correlated with the Y component of the IMF: when the magnitude of B_y was large compared to B_z , the rays appeared; otherwise, the distribution was relatively unstructured. Late in the flight the field of view of the imager included the field of view of flow measurements from the European incoherent scatter radar (EISCAT). The rays visible in 630.0-nm emission exactly aligned with the position of strong flow jets observed by EISCAT. We attribute this correspondence to the requirement of quasi-neutrality; namely, the soft electrons have their largest precipitating fluxes where the bulk of the ions precipitate. The ions, in regions of strong convective flow, are spread out farther along the flow path than in regions of weaker flow. The occurrence and direction of these flow bursts are controlled by the IMF in a manner consistent with newly opened flux tubes; i.e., when $|B_y| > |B_z|$, tension in the reconnected field lines produce east-west flow regions downstream of the ionospheric projection of the x line. We interpret the optical rays (flow bursts), which typically last between 5 and 15 min, as evidence of periods of enhanced dayside (or lobe) reconnection when $|B_y| > |B_z|$. The length of the reconnection pulse is difficult to determine, however, since strong zonal flows would be expected to persist until the tension force in the field line has decayed, even if the duration of the enhanced reconnection was relatively short.

1. Introduction

In an open magnetosphere, the cusps are funnel-shaped regions of open field lines. Particles which gain entry onto these field lines through magnetic reconnection at the dayside magnetopause may therefore precipitate directly into the ionosphere. Cusp field lines are magnetically connected to the ionosphere in a localized region near 12 MLT; as these field lines convect antisunward over the polar caps, they form the high-latitude boundary layer (mantle). Extending in local time on either side of the cusp is the cleft, whose (presumably closed) field lines map to the low-latitude boundary layer

(LLBL) along the flanks of the magnetosphere [Reiff, 1979; Heikkila, 1985; Newell and Meng, 1988]. The open field line portion of the LLBL (mapping to the low-altitude cleft equatorward of the cusp) consists of such recently merged field lines that only the highest-energy magnetosheath plasma has had time to precipitate [Reiff *et al.*, 1977; Onsager *et al.*, 1993; Lyons *et al.*, 1994].

The primary method of distinguishing the ionospheric projections of the cusp, mantle, and LLBL is through the classification of low-altitude particle precipitation, although plasma flow, optical, and magnetic characteristics may also be used. Ions injected at the dayside magnetopause undergo a time-of-flight dispersion as they travel to the ionosphere, resulting in a precipitation signature which decreases in energy with increasing latitude for southward interplanetary magnetic field (IMF) conditions [Shelley *et al.*, 1976; Reiff *et al.*, 1977]. The precipitating fluxes are largest slightly downstream of the last closed field line, where the ions with the highest phase space density (near 1 keV) precipitate [Onsager *et al.*, 1993]. As injected ions enter the tail lobe, the supply of plasma to low altitudes is cut off because the magnetosheath source flow becomes super-Alfvénic and tailward [Hill and Reiff, 1977; Onsager *et al.*, 1993]. Based on thousands of passes of DMSP data, Newell and Meng [1988, 1989] classify the low-altitude cusp as a region of intense, structureless, low-average-energy electron and ion precipitation, while cleft (LLBL) precipitation is characterized

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by more structure, lower flux, and higher average energy. The ionospheric footprint of the mantle, poleward of the cusp proper, is identified as the low-energy tail of the dispersed precipitation [Newell *et al.*, 1991]. Midaltitude measurements from DE 1 showed that these low-energy precipitating ions are associated with upgoing higher-energy ions which have mirrored at lower latitudes and thus will form the inner part of the plasma mantle [Burch *et al.*, 1982]. Observations of a reversed ion dispersion signature, indicative of sunward convection from lobe reconnection, have been observed for $B_z > 0$ [Burch *et al.*, 1980, 1986; Escoubet and Bosqued, 1989; Woch and Lundin, 1992; Basinska *et al.*, 1992]. At midaltitudes, the general dispersion signatures are also observed as a series of V's due to the pitch angle dependence of the field-aligned ion transit time [Burch *et al.*, 1982; Menietti and Burch, 1988].

The locus of cusp precipitation and the direction of cusp-region plasma flow are strongly dependent on the direction of the IMF. For instance, particle measurements have shown that cusp precipitation shifts (1) equatorward during intervals of $B_z < 0$ [e.g., Burch, 1973; Reiff *et al.*, 1980; Carbary and Meng, 1986; Newell *et al.*, 1989], (2) toward dawn for $B_y < 0$ ($B_y > 0$) in the northern (southern) hemisphere, and (3) toward dusk (or remaining near noon) for $B_y > 0$ ($B_y < 0$) in the northern (southern) hemisphere [Candidi *et al.*, 1983; Newell *et al.*, 1989; Aparicio *et al.*, 1991]. (The B_y shift is more apparent for $B_z < 0$ than $B_z > 0$). The direction of cusp-region plasma flow is predominately (1) antisunward (sunward) for $B_z < 0$ ($B_z > 0$), (2) toward dawn for $B_y > 0$ ($B_y < 0$) in the northern (southern) hemisphere, and (3) toward dusk for $B_y < 0$ ($B_y > 0$) in the northern (southern) hemisphere [e.g., Reiff and Burch, 1985; Heppner and Maynard, 1987]. The convection throat is located slightly on the prenoon side (in the corotating frame) regardless of the sign of B_y for a gardenhose IMF orientation [Heelis, 1984], while in the inertial frame the convection is more symmetric around noon [Maynard *et al.*, 1991b]. High time resolution radar measurements of flow direction on newly opened field lines (equatorward edge of the cusp for southward IMF) indicate that the east-west flow direction in the cusp responds very rapidly (of the order of a few minutes) to a reversal in IMF B_y [Clauer and Banks, 1986; Etemadi *et al.*, 1988; Greenwald *et al.*, 1990]. Observations of plasma flow reversals at low magnetopause latitudes also clearly support B_y dependent asymmetric polar cap convection. The B_y effect is attributed to tension in the newly opened field lines which alters the direction of ionospheric plasma flow [Cowley *et al.*, 1983, 1991; Reiff and Burch, 1985; Saunders, 1988; Gosling *et al.*, 1990; Maynard *et al.*, 1991a].

The question of whether the cusp arises from quasi-steady or transient reconnection is an important ongoing controversy. Although pulsed reconnection can yield a persistent cusp feature (assuming the repetition rate is faster than the ion precipitation time), it will create a discontinuous energy / latitude dispersion signature. A continuous energy / latitude dispersion signature, on the other hand, indicates quasi-stationary reconnection, injection, and convection. As we will show below, even continuous reconnection can cause a discontinuous dispersion signature, if the convective flow has a substantial, variable, component across the satellite trajectory [Reiff, 1990]. There is ample evidence for periods of apparently continuous reconnection at the magnetopause [e.g., Sonnerup *et al.*, 1981; Gosling *et al.*, 1982, 1986; Paschmann *et al.*, 1986, 1989] and for correspondingly quasi-

steady cusp ion dispersion signatures [e.g., Reiff *et al.*, 1977; Hill and Reiff, 1977; Frahm, 1984; Aparicio *et al.*, 1991; Mukai *et al.*, 1991; Phillips *et al.*, 1993]. There is also evidence supporting transient dayside merging, or flux transfer events (FTEs) [e.g., Russell and Elphic, 1979; Saunders *et al.*, 1984], their ionospheric flow jet signatures [Lockwood *et al.*, 1989a; Elphic *et al.*, 1990], and the correspondingly impulsive cusp ion dispersion signatures [e.g., Saflekos *et al.*, 1990; Bosqued *et al.*, 1991; Escoubet *et al.*, 1992; Lockwood and Smith, 1992; Yamauchi and Lundin, 1994]. It may also be possible that ionospheric flow bursts are caused by transients in the sheath magnetic field rather than in the reconnection rate [Newell and Sibeck, 1993], although Lockwood *et al.* [1994] dispute that possibility. In any case, it is likely that the cusp exhibits spatial or temporal enhancements superimposed on a background, steady state cusp driven by quasi-steady state reconnection on the dayside, at least for southward IMF [Smith and Lockwood, 1990; Southwood *et al.*, 1988; Scholer, 1988].

Recently, increased emphasis has been placed on studying magnetopause and boundary layer processes through the use of low-altitude and ground-based observations. Because of the requirement of quasi-neutrality, magnetosheath electrons precipitate with the injected ions, exciting 630.0 nm emissions and, to a lesser degree, emissions at 557.7 nm, 427.8 nm, and 391.4 nm. These emissions can be monitored by appropriately positioned photometers and all-sky cameras [e.g., Sandholt *et al.*, 1985; 1986a, b; Oguti *et al.*, 1988; Sandholt and Newell, 1992]. Likewise, density enhancements and ion flow irregularities in the dayside ionosphere may be probed using radars, which can provide high time resolution information on the density, temperature, and flow characteristics of the cusp and cleft regions [Goertz *et al.*, 1985; Todd *et al.*, 1986; Greenwald *et al.*, 1990; Lockwood *et al.*, 1989b, 1993; Watermann *et al.*, 1992; Baker *et al.*, 1990, this issue]. The Geospace Environment Modeling (GEM) Pilot program, on January 16, 1990 [Carlson and Basu, 1990], was designed as an intercalibration experiment to acquire joint ground-based, aircraft, and satellite data for the purpose of identifying the ionospheric signatures of the magnetospheric cusp and cleft. The Phillips Laboratory's Airborne Ionospheric Observatory (AIO) flew a course over the northern polar cap that kept it between 75° and 80° invariant latitude and within an hour of 12 MLT for over 5 hours. The AIO collected all-sky image, sounder, spectrometer, and scintillation data of the magnetospheric cusp during the flight. During the latter part of the flight, the field of view of the all-sky imager overlapped the field of view of the EISCAT radar. This paper will report the results of the coordinated optical, radar, and IMF data during the GEM Pilot program.

2. Observations

The flight path of the AIO, DE 1 satellite overflight, and the EISCAT field-of-view are shown in polar geographic coordinates in Figure 1. The 8-hour flight originated at Eielson Air Force Base, Alaska and ended at Andoya, Norway. From 0400 to 0645 UT the AIO, while keeping the dayside aurora in its field of view, remained within ± 15 min of noon magnetic local time. From 0645 to 0800 UT it drifted into the prenoon sector, crossing local times ranging from 1100 to 1040 MLT. IMF data, shown in Figure 2, are available from

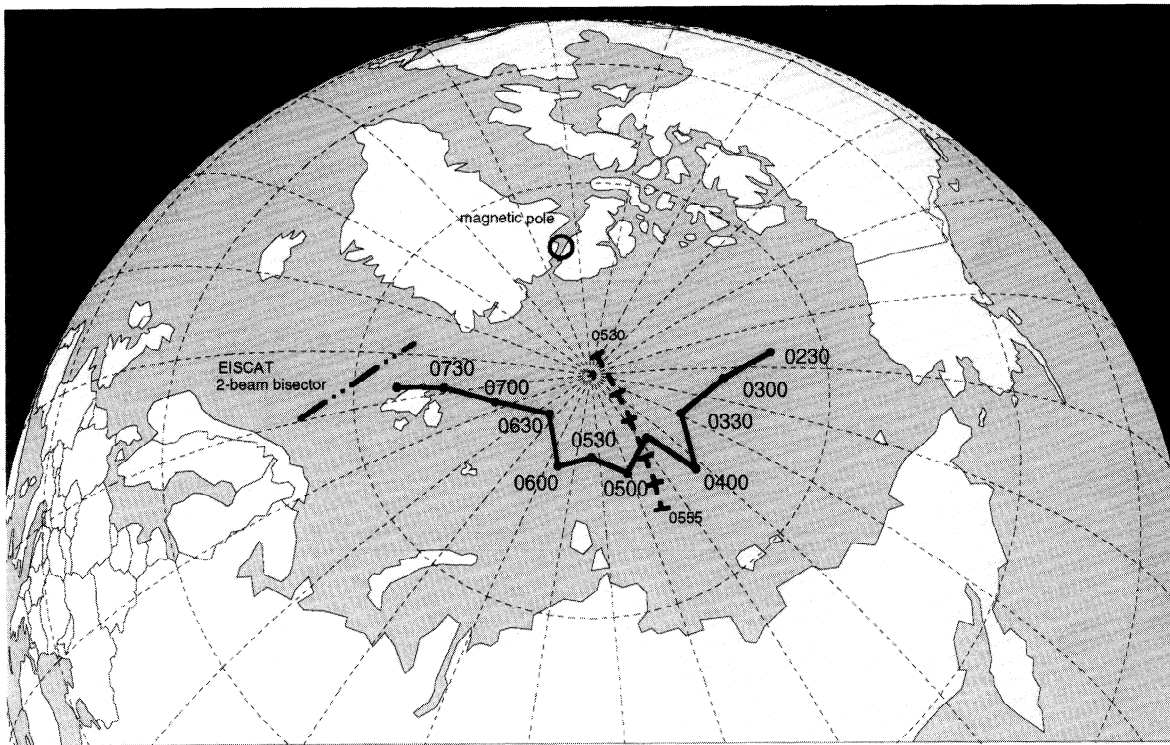


Figure 1. Polar geographic projection of the flight path of the AIO aircraft between 0230 and 0800 UT. The ground track of the DE 1 spacecraft (mapped to 250 km) is shown between 0530 and 0555 UT. From 0400 to 0645 UT the AIO adjusted its course to remain close to noon MLT while staying under the region of strongest 630.0-nm emission. The aircraft intercepted the EISCAT field of view near 0800 UT.

0530 UT onward. The B_z component was weakly northward during this period and B_y was positive, becoming negative only briefly from ~0640 to 0655 UT. The K_p 3-hour indices (3-, 3, 2) for the first half of January 16 indicate relatively quiet magnetospheric conditions.

During the flight the all-sky imaging photometer (ASIP) recorded 160° field-of-view images of the dayside aurora in three different wavelengths (427.8 N_2^+ , 557.7 $O[I]$, and 630.0 $O[I]$ nm). The three filters were cycled once per minute, with individual exposures requiring 2 s. In addition, two Ebert-Fastie scanning spectrometers measured the absolute intensities of the 391.4 nm N_2^+ ($I_{391.4} = 3 \times I_{427.8}$) and 630.0-nm emissions with a $5^\circ \times 7^\circ$ field of view centered in the aircraft zenith. The all-sky images were used for assessing the overall structure and dynamics of the cusp-region aurora, while the spectrometer provided accurately calibrated absolute intensities (and intensity ratios). The 630.0-nm emission intensities (~2–4 kR) and the $I(630.0 \text{ nm}) / I(391.4 \text{ nm})$ ratios (~10–15) are consistent with the soft (< 200 eV) electron precipitation expected for cusp aurora [Rees and Luckey, 1974; Cogger *et al.*, 1977; Eather *et al.*, 1979; Sandholt, 1988].

The most striking morphological change in the ASIP images was the appearance of zonally elongated fingers of emission, or "rays," whenever the magnitude of B_y was large compared to the magnitude of B_z . The four all-sky images shown in Figure 3 illustrate this correlation. These images have been transformed to geographic coordinates (north at the top) using an assumed emission height of 250 km. The magnitudes of B_y and B_z are noted in the upper right of each image. The images at 0535 and 0745 UT display elongated

emission features stretching (predominately) to the west, while the other two images show relatively unstructured emission.

The 1-min resolution images were closely examined during the appearance of the elongated features. For the event at 0535 UT, they clearly show that two distinct faint arcs brightened to the west of the AIO to form the features shown in Figure 3. After approximately 6–10 min, the bright arcs faded and the coherent structure to the west dissipated. Since the exposure time of each image is 2 s, there is very little blurring during the exposure. The intensity of the emissions are slightly diminished at the edges of the images, however, due to a combination of reduced sensitivity and lens defocusing. This vignetting can cause the 630.0-nm aurora to appear localized, when it may actually extend beyond the field of view. Since all images are subject to the same vignetting, however, we believe that differences in auroral structure between successive images are real.

The DE 1 satellite, traveling from pole to equator, passed through the ASIP field of view between 0530 and 0555 UT. Precipitating H^+ and O^+ ions, measured by the energetic ion composition spectrometer (EICS) are shown in Plate 1. Sheathlike H^+ ions are encountered between 0545 and 0549 UT. The energies of the H^+ ions display a slight reverse dispersion, decreasing in energy from high to low latitudes (0547 – 0549 UT). The region of most intense H^+ precipitation, marked by a box on the footprint of the DE trajectory in Figure 3, occurs at the same invariant latitude as the maximum intensity 630.0-nm emissions. The reverse dispersion is also seen in the weak, precipitating O^+ fluxes.

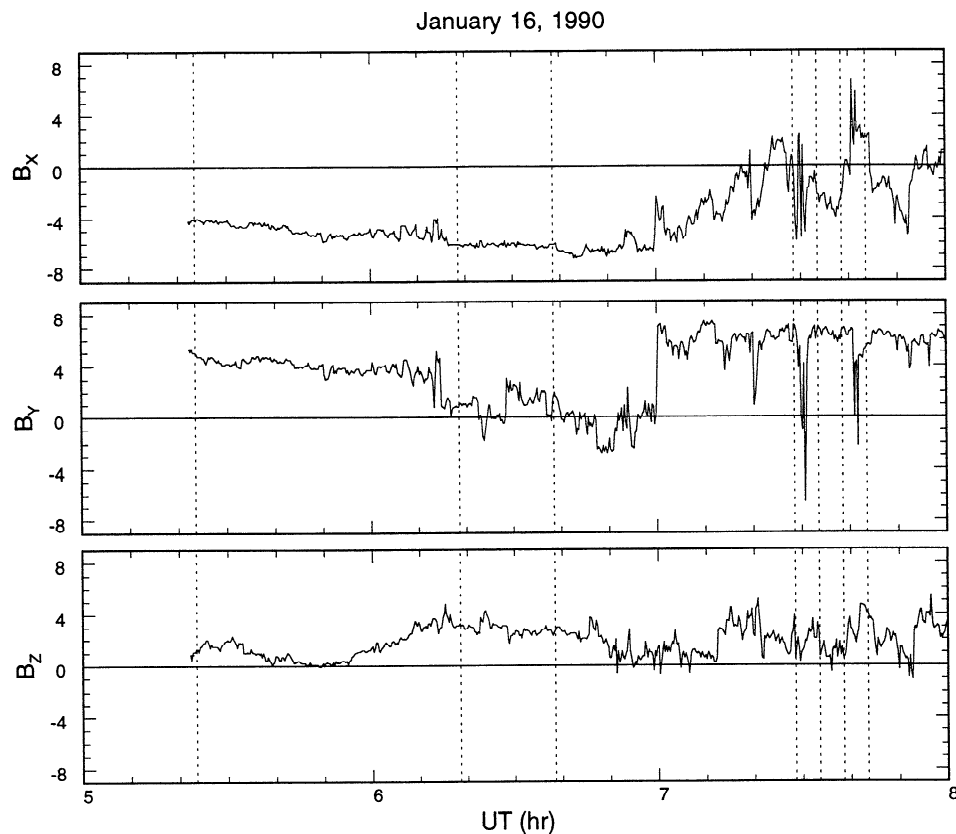


Figure 2. GSM B_x , B_y , and B_z components of the IMF for the period 0500 - 0800 UT, January 16, 1990. The IMP 8 spacecraft was located at an average GSM position of (29.3, 11.7, 7.4) R_E , corresponding to a solar wind delay time of about 11.5 min. The vertical dashed lines indicate the IMF conditions for the events in Figures 3 and 5. (High time resolution data courtesy Ron Lepping and NDADS.)

The source of this O^+ may be ions, originally part of the cleft ion fountain or turbulent entry layer, which are accelerated earthward as a result of lobe merging near the cusp.

Toward the end of the flight the ASIP and EISCAT fields of view overlapped. The EISCAT radar was using the "Polar" beamswing technique to obtain two-dimensional ionospheric flow vectors in the region north of Tromsø. The technique is described in detail by Willis *et al.* [1986] and Lockwood [1991]. Figure 4 shows the derived (at the bisector of the two radar beams) velocity vectors overlaid on a series of four images from 0740 to 0755 UT. During this time, $B_y > B_z > 0$ (see Figure 2). The flow vectors clearly show that the 630.0-nm rays are aligned with latitudinally narrow, strong flows in the ion velocity (~ 1 km/s), which persist for 5-15 min. The strong westward flow is consistent with the field curvature force on newly opened field lines for the observed positive IMF B_y . If these flow bursts are indicative of newly opened field lines as we believe, the cusp footprint extends at least to 1040 MLT (the location of the AIO) at this time.

3. Discussion

The fact that the high-flow regions in the EISCAT data coincide with elongated emission features can be explained by the mechanism illustrated in Figure 5. A newly reconnected field line will be transported into the tail lobe, first under the action of the field line curvature (tension) force and

subsequently by magnetosheath flow. There is a limited time dt for which the precipitation yielding strong 630.0-nm emission is seen on each newly opened (or reconfigured) field line, due to the reduction of the injected fluxes near the magnetosheath Alfvén line. (Recall that quasi-neutrality allows strong magnetosheath electron precipitation only where the bulk of the ions precipitate [Burch, 1985]). This time dt will depend upon the reconnection site, the sheath flow, and the sheath field orientation. Under high-flow conditions ions injected at the merging point experience a greater degree of velocity dispersion, and thus a longer ground track, than ions injected under slow-flow conditions (Figure 5a). If strong- B_y conditions also exist, the ion ground track and 630.0-nm emission will be azimuthally extended (in the direction opposite to B_y) downstream of the x line (Figures 5b and 5c). Even lobe reconnection (Figure 5c) yields flows extending to the dawnside for B_y positive conditions [Crooker, 1979, Reiff and Burch, 1985].

With this explanation in mind, we interpret the periodic appearance of azimuthally extended emission features in the earlier all-sky images (those without coincident EISCAT measurements) as indirect evidence of enhanced ionospheric flow (i.e., periods of enhanced dayside or lobe reconnection). The occurrence and direction of the optical rays / flow bursts are controlled by the IMF in a manner consistent with the tension forces on newly opened flux tubes: in the first and fourth images in Figure 3, rays appear to emanate to the west

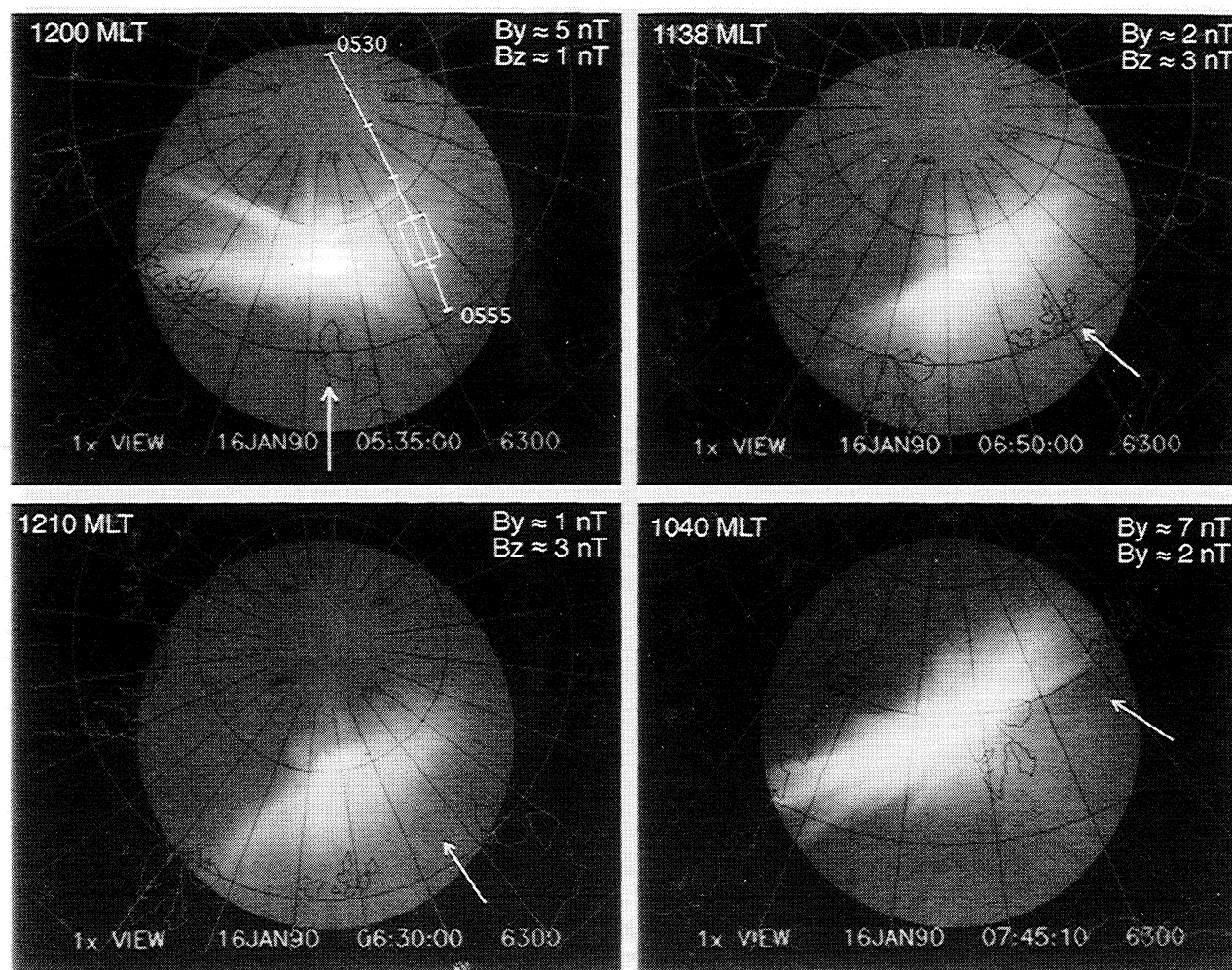


Figure 3. Four representative 630.0-nm all-sky images illustrating the correlation between the appearance of zonally elongated fingers of emission and the ratio of $|B_y|/|B_z|$. When the B_y component is larger than B_z (top left and bottom right), rays emanate from the centroid of the precipitation; otherwise the distribution is relatively unstructured. The local time of each image is noted at the top left and the 12 MLT meridian is marked with an arrow.

from a basically oval centroid, consistent with the dawn-directed convection velocity expected for $B_y \gg 0$ [e.g., Crooker, 1979, 1988; Reiff and Burch, 1985]. In the other two images, $|B_y| \leq |B_z|$, so highly asymmetric dawn-dusk flows are not expected (and are not observed). During the short period of time when $B_y < 0$ (0640 - 0655 UT), one might have expected the location of the cusp to shift downward and to appear with dusk-directed flow jets; perhaps due to the short duration of negative B_y , neither signature appears in the images (but could have occurred out of our field of view to the east). We also suggest that while the central portion of the 630.0-nm emission represents the "cusp proper" where the 1-2 keV ions precipitate, the extended rays are the footprint of the low-energy extension of the cusp precipitation (i.e., the mantle). Pinnock *et al.* [1993] also showed DMSP and PACE HF radar observations which clearly demonstrated an extended, enhanced convection channel downstream of the cusp ion precipitation region (cf. their Plates 2 and 3).

The dispersion signature observed by low-altitude satellites passing through the cusp depends on the satellite's proximity to the merging line(s) and on any discontinuities in either the

merging line or the ionospheric flow. Figures 5d, 5e, and 5f schematically illustrate these dependencies by showing for each case (slow flow, strong azimuthal flow, and strong azimuthal flow with lobe reconnection) the precipitating ion energies expected along two sample flow lines. Also shown are possible crossings of the region by polar-orbiting spacecraft (dotted lines). As discussed above, an elongated dispersion is observed in regions where the ion flow velocity is largest (e.g. the poleward flow paths in Figures 5e and 5f). Note in Figure 5d that almost any crossing of the distribution would show a "normal" energy dispersion, i.e., energies falling from low to high latitudes. In Figures 5e and 5f, however, the observed dispersion depends on how the spacecraft trajectory cuts through the region. In Figure 5e, a spacecraft traveling from point *q* to point *r* would see a dispersion that falls in latitude until it reached the high-flow region, where the energy would rise before falling again. If the dayside *x* line were broken into segments or if there were a truly discontinuous flow velocity in the ionosphere (caused by a flow discontinuity in the magnetosheath), the dispersion signature would appear to jump up discontinuously. Thus a

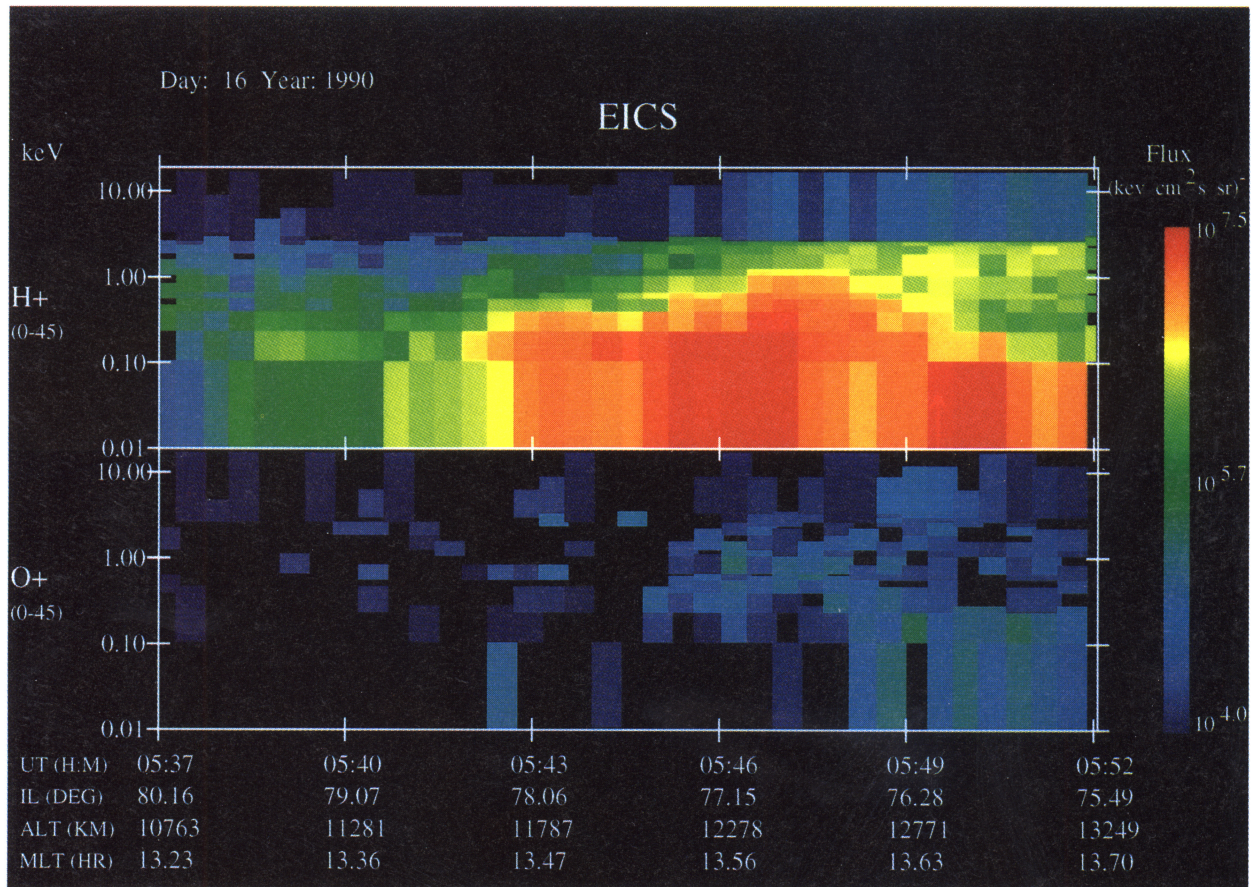


Plate 1. Precipitating H^+ and O^+ ions measured by DE 1 during its passage through the northern hemisphere cusp region. The satellite was traveling equatorward at an altitude of about 12,000 km. The most intense fluxes of magnetosheath H^+ ions ($\sim 10^7$ keV cm $^{-2}$ s $^{-1}$ sr $^{-1}$) are observed between 0545 and 0549 UT. Reversed dispersion is seen in both H^+ and O^+ between 0547 and 0549 UT, with peak O^+ fluxes $\sim 10^5$ keV cm $^{-2}$ s $^{-1}$ sr $^{-1}$.

"stair-step" dispersion signature as seen by *Escoubet et al.* [1992] may not necessarily be a result of impulsive reconnection, but could simply represent a discontinuity in the flow velocity across the spacecraft trajectory. This latter interpretation is supported by recent observations of a stairstep dispersion signature recorded by DE 1 and 2 at nearly the same local time over 20 min apart [*Onsager et al.*, 1995].

Figure 5f shows the ion dispersion energies expected during periods of lobe reconnection. A spacecraft crossing from u to v would see a truly discontinuous ion dispersion as it crossed from a flow path connected to the dayside x line to a flow path connected to the lobe x line. Whether that discontinuity would be a step up or down in energy would depend on the difference in the time along the flow path from each x line. A spacecraft crossing from y to z would see a "reverse" dispersion signature: highest energies at the highest latitudes, much as was seen in this case by DE 1. A spacecraft crossing from x to z would see highest energies at both high and low latitudes, with a minimum energy in between. We suggest that this type of cusp configuration is the source of the large-scale ion V distributions observed for northward IMF [*Reiff et al.*, 1977, 1980; *Burch et al.*, 1980]. Similar large-scale V distributions can be seen in the cusp data of *Woch and Lundin* [1992; Figure 3] and *Yamauchi et al.* [1993, Figure 2]; the former also shows large zonal flows at the edges of the cusp precipitation region.

Figure 6 shows an all-sky image taken at 0545 UT, close to the time that DE 1 saw reversed-dispersion cusp ion precipitation. This dispersion signature indicates a sunward convection component in the cusp region at this time, or at least a satellite trajectory in which the poleward portion of the trajectory passed closest to the lobe x line (e.g., path y - z in Figure 5f). We have overlaid a possible convection pattern (similar to Figure 5c) which accounts for a reverse dispersion within a generally westward flow. One possible location of the x line consistent with the observed dispersion is shown with dotted lines: one segment corresponds to magnetic merging on the dayside, and the other segment (nearer DE 1) corresponds to magnetic merging of the IMF with lobe field lines. Merging in either location will yield a westward outflow consistent with the rays, but the proximity of the spacecraft trajectory to the lobe x line is more consistent with the reversed energy dispersion observed by DE 1.

Finally, we note that when the flow-aligned jets appear in the images, they typically persist for 5-15 min, with one event lasting 25 min. *Lockwood and Cowley* [1988] and *Cowley and Lockwood* [1992] make the point that once a field line has been opened (or reconfigured by lobe reconnection), zonal flows in the ionosphere will last for an extended period while the curvature force decays (i.e., the field line straightens). If the flow jets (extended emission features) are the result of enhanced reconnection pulses when $|B_y| > |B_z|$,

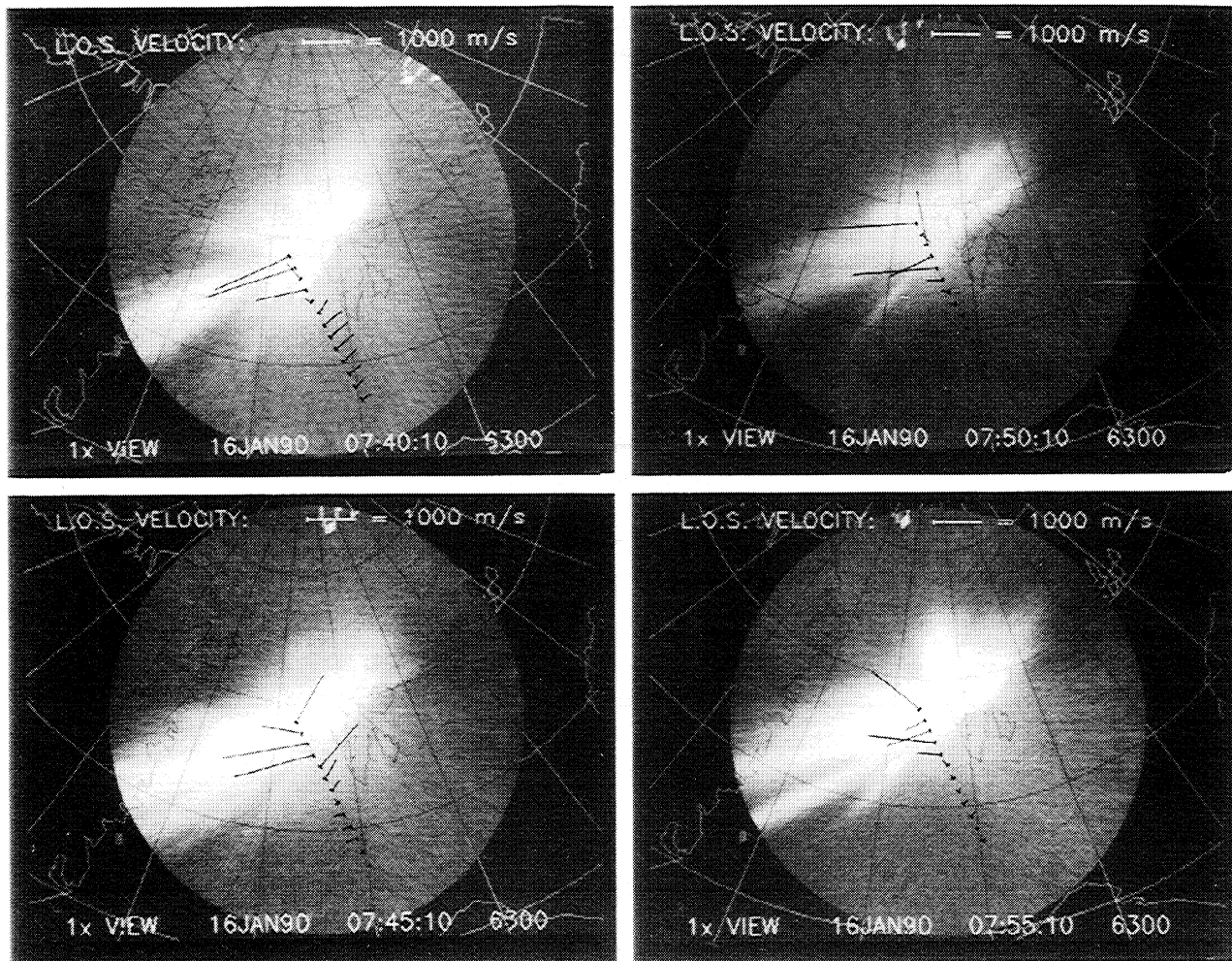


Figure 4. Two-dimensional flow velocities, derived from the "Polar" beamswing technique, superimposed on four ASIP images from 0740 to 0755. The elongated optical emissions are aligned along the direction of the high flow velocities.

the duration of the pulses must be less than the observed durations of the optical rays (if the pulses are sufficiently close together in time, the rays will appear continuous, which they do not). The extent of the 630.0-nm emission along a flow stream line is $V dt$, where V is the convection speed and dt is the duration of the precipitation on any one newly opened field line which gives rise to 630.0-nm emission. The overlapped optical / EISCAT observations reveal that the extent of the rays ($V dt$) and the convection speed (V) are roughly correlated and hence that the interval dt is relatively constant during this time. Given that the optical jets extend at least 500 km and the EISCAT data reveal that within them V is roughly 1 km/s, we derive a minimum value of dt of 500 s. This is quite reasonable; a 150–200 eV ion (typically the lowest energy cusp ions for which the electron precipitation is significant) has a field-aligned velocity of 170–200 km/s; 500 s corresponds therefore to a field-aligned distance of 13–15 R_E to the x line, which is reasonable given magnetic field models and is consistent with previous estimates [e.g., Carlson and Torbert, 1980; Menietti and Burch, 1988].

4. Summary

The observations presented here show transient dayside auroral jets, associated with plasma flow enhancements, which

behave in the same manner as that suggested for newly opened flux tubes; i.e., flow jets appear when $|B_y| > B_z$ and are directed opposite to the direction of B_y . A simple model of why the high-flow regions appear as elongated optical rays is suggested: under high-flow conditions (i.e., enhanced reconnection electric fields) ions injected at the merging point have a longer ground track $V dt$ than those injected under slow-flow (background) conditions; due to the requirement of quasi-neutrality, the region of precipitating electron flux and thus 630.0-nm emission is also longer. It is significant that the flow jets are intermittent and transient, with durations of tens of minutes. The optical jets are interpreted as evidence of enhanced reconnection which, under northward IMF conditions, occurs predominantly when $|B_y| > B_z$. We have presented a possible convection pattern, involving both dayside and lobe reconnection, which accounts for the reversed ion dispersion seen by DE 1 in conjunction with the strong westerly flow. It is suggested that if periods of enhanced reconnection occur at these locations, the duration of the reconnection is less than the duration of the flow-aligned jets, which varies between 5 and 25 min. Given that the precipitation within the jets lasts for ≥ 8 min, we conclude that the reconnection pulses last for ~ 0 –17 min. One must keep in mind, however, that little direct evidence exists which links the transient dayside ionospheric phenomena discussed

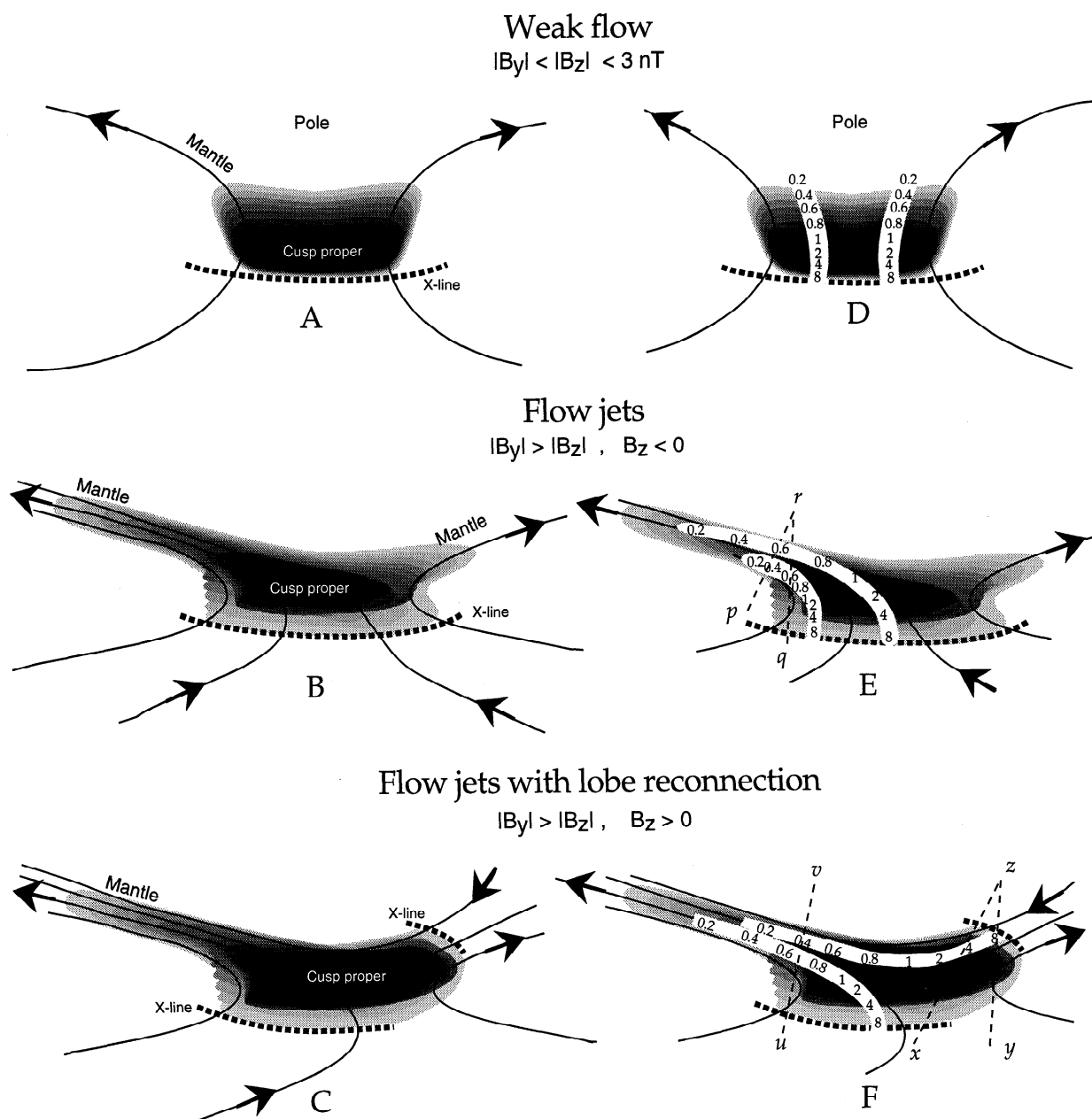


Figure 5. Schematic diagram showing cusp precipitation patterns (and 6300 \AA emission) expected for times of (a) weak flow, and (b and c) strong B_y -dependent flow. Noon is toward the bottom of each figure. (c) differs from Figure 5b in that some of the westward directed flow originates from a lobe merging x line in the polar cap. The flux is most intense (darkest) slightly downstream from the x line, where the 1-2 keV ions precipitate. The fluxes are also largest near noon because those field lines map to the stagnation region at the nose of the magnetosphere, where the sheath density is the largest. Flow paths which cross the x line farther from noon have weaker source densities and thus lower fluxes. The contours are stretched in the direction of the flow, since the ion dispersion distance $v dt$ is larger where the convection velocity is larger (dt is observed to be $\sim 500 \text{ s}$). Sample precipitating ion energies (keV) for each of the flow configurations are shown on the left. Lobe reconnection (f) can yield a reversed dispersion signature on the afternoon side as observed by DE 1, which traveled approximately from point z to point y. A spacecraft traveling from x to z would observe a V-shaped ion dispersion, with the minimum energy at the transition between sources.

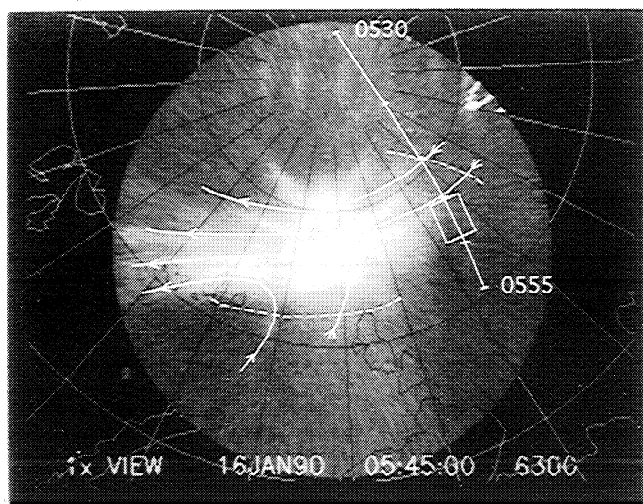


Figure 6. Possible ionospheric flow lines (similar to Figure 5f) overlaid on an all-sky image at 0545 UT. The flow paths are consistent with the optical jets and the reverse dispersion observed by EICS on DE 1.

here to FTE signatures at the magnetopause; the only examples being those presented by *Elphic et al.* [1990] and *Sandholt et al.* [1992].

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